Sub-Fourier sensitivity in ac driven quantum systems

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Avoided crossing and sub-Fourier sensitivity in ac-driven quantum systems

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Model	Intro	q-ratchet	Avoided crossings	Theory	Demonstration
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Outline					

- 1 Model system
- 2 What is sub-Fourier Sensing?
- 3 Quantum ratchet
- 4 The quantum ratchet: Avoided crossings
- 5 Exploiting avoided crossings: the theory
- 6 Exploiting avoided crossings: implementation

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• A single particle x(t) (quantum or classical)

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Ratchet:

$$\langle v \rangle = \lim_{t \to \infty} \frac{\langle x(t) \rangle - \langle x(0) \rangle}{t}$$

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$$f(t) = \int_{-\infty}^{\infty} d\omega \, F(\omega) e^{i\omega t}$$

$$e^{i\omega t} = \cos(\omega t) + i\sin(\omega t)$$

 $\omega = \frac{2\pi}{T}$



Joseph Fourier

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$$f(t) = \int_{-\infty}^{\infty} d\omega F(\omega) e^{i\omega t}$$

• Fourier inequality

 $\Delta \omega \Delta t \geq 2\pi$

 $\Delta \omega$: width of $F(\omega)$

 Δt : width of f(t)



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A sensor based on resonances:

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• For every periodic driving F(t),

$$F(t+T) = F(t), \quad \omega = \frac{2\pi}{T}$$

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$$\langle v \rangle = \lim_{T_s \to \infty} \frac{1}{T_s} \int_0^{T_s} dt \, v(t)$$

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• Fourier limit: (for linear systems) System bandwidth $\Delta \omega \sim \frac{2\pi}{T_s}$ T_s : interaction or observation time

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$$\Delta\omega\ll\frac{2\pi}{T_s}$$



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 Example: guantum $\delta\text{-kicked-rotor}$

David Cubero Brownian ratchets

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Cold atoms exposed to N pulses of off-resonant standing waves of light

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FIG. 1. Below the Fourier limit. (a) Experimental measurement of the zero-velocity atom number, p(0), as a function of the ratio $r = f_2/f_1$ of the two excitation frequencies. Parameters: $f_1 = 18$ kHz, K = 42, $N_1 = 10$, $\tau = 3 \mu$ s and in order to avoid pulses overlap we set $\varphi = \pi$. Averaging: 100 times. (b) $F_{1/2}(r)$, for comparison with the Fourier transform of the kick sequence (amplitude and offset are arbitrary for $F_{1/2}$).

$$\Delta f_2 T \approx \frac{1}{38} \ll 1. \tag{3}$$



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FIG. 1. Experimentally measured fidelity distribution due to five kicks of strength $\phi_d = 0.8$ followed by a shifted kick of strength $5\phi_d$. The mean energy (triangle same five kick rotor is shown for comparison. Numeric lations of the experiment for a condensate with momentu $0.06\hbar G$ are also plotted for fidelity (dashed line) ar energy (solid line). The amplitude and offset of the s

• Sub-Fourier sensing: (non-linear systems)

$$\Delta\omega\ll\frac{2\pi}{T_s}$$

A second example:

classical rocked ratchet with biharmonic driving

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A second example:

classical rocked ratchet with biharmonic driving

$$F(t) = F_1 \cos(\omega_1 t) + F_2 \cos(\omega_2 t + \theta)$$

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Resonances at $\omega_2 = (p/q)\omega_1$
SubFourier sensitivity:

$$\Delta\omega_2 = \frac{2\pi}{q T_s}$$

Casado-Pascual, DC & Renzoni, PRE 88, 062919 (2013). DC, Casado-Pascual, & Renzoni, PRL 112, 174102 (2014).

What is sub-Fourier sensing?

• Sub-Fourier sensing: (non-linear systems)

 $\Delta\omega\ll\frac{2\pi}{\tau}$

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FIG. 1 (color online). Current vs driving frequency ω_2 for the overdamped system (1) with the driving F_2 and $F_0 = 5.75$. Reduced units are defined in the simulations such that $U_0 = k = \gamma = \omega_1 = 1$. Empty and filled diamonds correspond to $T_s = 10^4$ and 10^5 , respectively. The lines are the predictions given by (6) with q = 113, p = 355, and $v_0 = 1/(2q)$. The horizontal bars depict the frequency width (7), showing a resolution 113 times smaller than that expected from the Fourier width $2\pi/T_s$.

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• No dissipation, $H = p^2/(2m) + V(x) - xF(t)$

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• Bloch-Floquet states with current $v_n = \frac{1}{\hbar} \frac{\partial \epsilon_n}{\partial k}$

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 $V(x) = V_0 \cos(2\pi x/L), \ F(t) = F_1 \cos(\omega_1 t) + F_2 \cos(\omega_2 t + \theta), \ \omega_2 = 2\omega_1, \ \theta = -\pi/2$

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The quantum ratchet: Avoided crossings



Denisov, Morales-Molina, Flach & Hänggi, Phys. Rev. A 75 063424 (2007)



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Frequency dependence near resonances?

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• Bloch-Floquet states:

$$v_n(k,\theta) = \frac{1}{\overline{T}} \int_{t_0}^{t_0+T} dt \langle \psi_{k,n}(t) | (p/m) | \psi_{k,n}(t) \rangle$$



Frequency dependence near resonances?

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- Finite-time current: $v_{T_s} = \frac{1}{T_s} \int_{t_0}^{t_0+T_s} dt v(t)$ $v(t) = \langle \psi(t) | (p/m) | \psi(t) \rangle$

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 Intro
 q-ratchet
 Avoided crossings
 Theory
 Demonstration

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- Finite-time current: $v_{T_s} = \frac{1}{T_s} \int_{t_0}^{t_0+T_s} dt v(t)$ $v(t) = \langle \psi(t) | (\rho/m) | \psi(t) \rangle$
- Asymptotic approximation: DC & Renzoni, PRE 97, 062139 (2018).

$$v_{T_s} \sim \frac{1}{\Delta \omega_2 T_s} \int_{\theta_0}^{\theta_0 + \Delta \omega_2 T_s} d\widetilde{\theta} \sum_{k_0, n} |C_{k_0, n}|^2 v_n \left(k(\widetilde{\theta}), \widetilde{\theta}\right),$$

 $\Delta\omega_{2} = \omega_{2} - \omega_{1}p/q, \quad \theta_{0} = \theta + \omega_{2}t_{0},$ $k(\widetilde{\theta}) = k_{0} + \lim_{\Delta\omega_{2} \to 0} \int_{t_{0}}^{t_{0} + T \lfloor \frac{\widetilde{\theta} - \theta}{\Delta\omega_{2}T} \rfloor} dt' F(t')/\hbar,$ $|C_{k_{0}, p}|^{2} = |\langle \psi_{k_{0}, p}(t_{0}) | \psi(t_{0}) \rangle|^{2}$

 Model
 Intro
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$$v_{T_s} \sim \frac{1}{\Delta \omega_2 T_s} \int_{\theta_0}^{\theta_0 + \Delta \omega_2 T_s} d\widetilde{\theta} \sum_{k_0, n} |C_{k_0, n}|^2 v_n \left(k(\widetilde{\theta}), \widetilde{\theta}\right),$$

$$\Delta\omega_2 = \omega_2 - \omega_1 p/q, \quad \theta_0 = \theta + \omega_2 t_0,$$

 $k(\theta) = k_0 + \lim_{\Delta \omega_2 \to 0} \int_{t_0}^{\infty + \tau + \tau \Delta \omega_2 \tau} dt' F(t')/\hbar_{t_0}$

 $|C_{k_0,n}|^2 = |\langle \psi_{k_0,n}(t_0)|\psi(t_0)\rangle|^2$

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 Intro
 q-ratchet
 Avoided crossings
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$$\Delta\omega_{2} = \omega_{2} - \omega_{1}p/q, \quad \theta_{0} = \theta + \omega_{2}t_{0},$$

$$k(\tilde{\theta}) = k_{0} + \lim_{\Delta\omega_{2} \to 0} \int_{t_{0}}^{t_{0} + T \lfloor \frac{\tilde{\theta} - \theta}{\Delta\omega_{2}T} \rfloor} dt' F(t')/\hbar,$$

 Model
 Intro
 q-ratchet
 Avoided crossings
 Theory
 Demonstration

 Composition
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Frequency dependence near resonances?

- Bloch-Floquet states: $v_n(k,\theta) = \frac{1}{T} \int_{t_0}^{t_0+T} dt \langle \psi_{k,n}(t) | (p/m) | \psi_{k,n}(t) \rangle$
- Finite-time current: $v_{T_s} = \frac{1}{T_s} \int_{t_0}^{t_0+T_s} dt v(t)$ $v(t) = \langle \psi(t) | (\rho/m) | \psi(t) \rangle$
- Asymptotic approximation: DC & Renzoni, PRE 97, 062139 (2018).

$$v_{T_s} \sim \frac{1}{\Delta \omega_2 T_s} \int_{\theta_0}^{\theta_0 + \Delta \omega_2 T_s} d\widetilde{\theta} \sum_{k_0, n} |C_{k_0, n}|^2 v_n \left(k(\widetilde{\theta}), \widetilde{\theta}\right),$$

$$\begin{split} \Delta \omega_2 &= \omega_2 - \omega_1 p/q, \quad \theta_0 = \theta + \omega_2 t_0, \\ k(\widetilde{\theta}) &= k_0 + \lim_{\Delta \omega_2 \to 0} \int_{t_0}^{t_0 + T \lfloor \frac{\widetilde{\theta} - \theta}{\Delta \omega_2 T} \rfloor} dt' F(t') / \hbar, \\ &|C_{k_0,n}|^2 = |\langle \psi_{k_0,n}(t_0) | \psi(t_0) \rangle|^2. \end{split}$$

Model	Intro	q-ratchet	Avoided crossings	Theory	Demonstration
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Outline					

- 1 Model system
- 2 What is sub-Fourier Sensing?
- 3 Quantum ratchet
- 4 The quantum ratchet: Avoided crossings
- 5 Exploiting avoided crossings: the theory
- 6 Exploiting avoided crossings: implementation

















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 $\psi(x, t = -2T) = \text{const.}, F_0 \text{ during } 2T \text{ such as to start from the right } k_0$. Top (+) has $\theta = -1.0851$, bottom (-) has $\theta = -1.0845$.





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 - Thank you for you attention!